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By D'Alembert's formula: $u(x, t) = 0 + \dots 2\pi \cdot (\sinh(\pi) - \sinh(-\pi)) = \dots \varphi(x) = -\lambda\varphi(x)$. $G(t) = (4 - \lambda)G(t)$. As noted in the hint, $\varphi(x) = \sin(nx)$, ...
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$\varphi(x) + \lambda\varphi(x) = 0$. $\varphi(0) = \varphi(L) = 0$. (1.6). and the equation: $G(t) = \lambda k G(t)$ $n=1$. [C. n. cos $(c/\lambda \cdot n \cdot t) + D \dots 2\pi \cdot 0 \cdot g(x)\varphi \cdot n \cdot (x)dx$. D ...
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$\lambda\Phi$. 4. theory. In particular, if one considers the Ginzburg-Landau model confined between two parallel ... $2(2\pi)$. $D/2 \cdot \infty$. $n=1$. $m(T, L)$. nL . $(D-2)/2$...
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After direct summation by the rule. ∞ . $n=1$. $a \cdot n \cdot n = -\ln(1 - a)$. (4.4). we have. $U \cdot sc = - \cdot 1$.
 2. $d \cdot 4 \cdot k \cdot (2\pi)$. $4 \cdot \ln(1 + (2m \cdot 2 + \lambda\varphi \dots$
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Holographic viewing device, and computer-generated hologram for ...

$D(u, v) = \lambda\Phi \text{HOLO}(u, v) / \{2\pi(n_1 - n_0)\}$ (1) Here λ is the center wavelength used, and n_1 and n_0 are the refractive indices of two materials that form the ...
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2 and 3 is etched to a depth of $\lambda\Phi \cdot \sup(1)m, n / 2\pi$ modulo λ , and right half ... case of equations (12) and (13), specifically, the case in which $p_m/n = 1/2$
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$d. n. f. dx. n. (-i\omega). n. F(\omega). (x\text{-derivatives}). fg = \int_{-\infty}^{\infty} f(x - \bar{x})g(\bar{x})d\bar{x} \dots d. 2. dx.$

2. $\varphi(x) = -\lambda\varphi(x)$ and Fourier Formulæ. Boundary ...

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$\Delta\varphi + \lambda\varphi = 0, \varphi(r, \theta, 0) = 0, \varphi(r, \theta, H) = 0, \partial\varphi/\partial r(a, \theta, z) = 0 \dots g(\theta + 2\pi) = g(\theta), \dots dg/d\theta(\theta$

$+ 2\pi) = dg/d\theta(\theta). r. 2. d \dots$

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$\psi(n) = P(n)\psi(n) = P(n)(\lambda\psi(n-1) + u(n)d(n)) = P(n)(\lambda\Phi(n-1)w(n-1) + u(n)d(n)). = P(n)((\Phi$

$(n) - u(n)u(n) \dots 2\pi. W. (n-2)).] , n = 1, 2, 3 \dots$

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